UBC Physics 102

Lecture 5 Version 2

Rik Blok

Outline

- ▷ Electric field
- ▷ Conductors
- ▷ Continuous charge distributions
- ▷ End

Electric field [Text: Sect. 21-6]

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  - Force depends on charge q but $\mathbf{E}$ is the same for *all test charges*.
  - So electric field is more useful quantity to work with. Once you know $\mathbf{E}$ can easily compute force on any test charge q via

$$\mathbf{F} = q\mathbf{E}.$$

Electric field, contd

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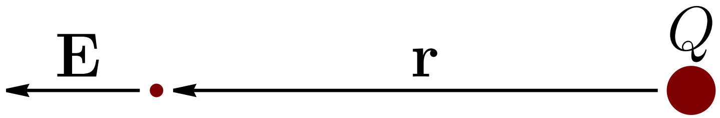
Electric field, contd

- **Definition:** *Coulomb's law*
    - Convenient to use electric field form of Coulomb's law.
    - Gives field at any point due to charge Q,



$$\mathbf{E} = k\frac{Q}{r^2}\hat{\mathbf{r}}.$$

Electric field, contd

- **Discussion: Superposition principle**

Electric field, contd

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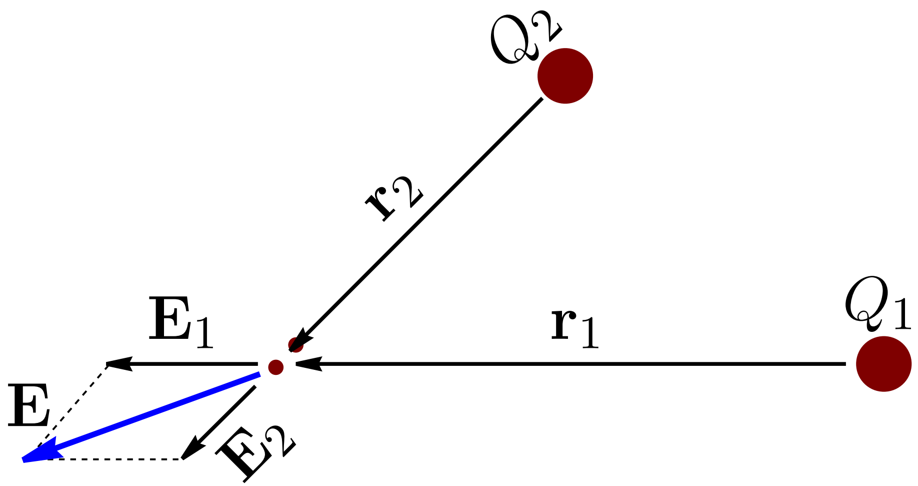
  - If dealing with more than one charge, can just add up electric field due to each to calculate net electric field at a point,



$$\mathbf{E} = \mathbf{E}_1 + \mathbf{E}_2 + \cdots .$$

Electric field, contd

- **Example: Pr. 40**

Electric field, contd

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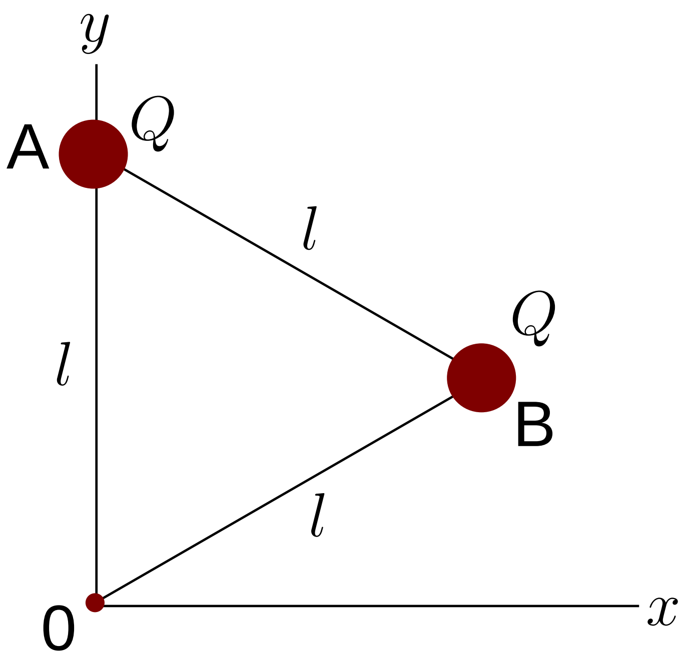
  - Determine the electric field $\mathbf{E}$ at the origin 0 due to the two charges at A and B.



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  - By the superposition principle

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  - Just need to find the directions. Direction from A to origin is

$$\hat{\mathbf{r}}_A = -\hat{\mathbf{j}}.$$

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  - B is at (x, y) where $y = \frac{1}{2}l$ and $x^2 + y^2 = l^2$ so $x = \sqrt{\frac{3}{4}}l$.

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  - Direction from B to origin is

$$\hat{\mathbf{r}}_B = -\sqrt{\frac{3}{4}}\hat{\mathbf{i}} - \frac{1}{2}\hat{\mathbf{j}} = -\frac{\sqrt{3}}{2}\hat{\mathbf{i}} - \frac{1}{2}\hat{\mathbf{j}}.$$

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  - So net electric field at origin is

$$\begin{aligned}\mathbf{E} &= \frac{kQ}{l^2}\left(\hat{\mathbf{r}}_A + \hat{\mathbf{r}}_B\right) \\ &= \frac{kQ}{l^2}\left(-\frac{\sqrt{3}}{2}\hat{\mathbf{i}} - \frac{3}{2}\hat{\mathbf{j}}\right). \quad \square\end{aligned}$$

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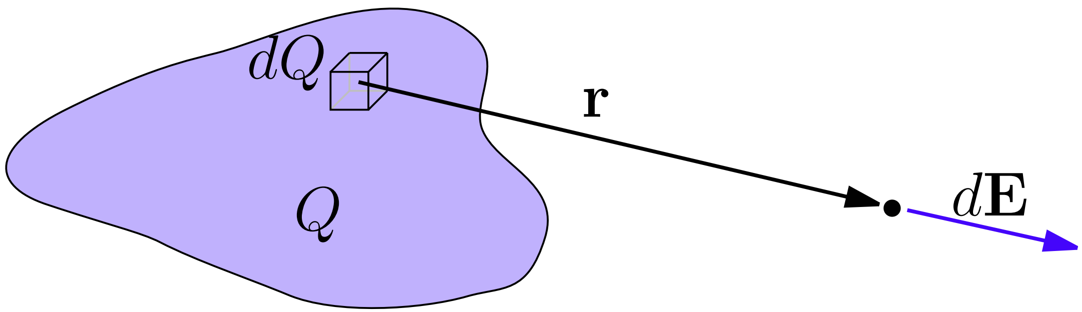
  - If object too large to be treated as point charge, can still solve for electric field.
  - Divide object into small chunks and add up field due to each chunk (superposition principle).



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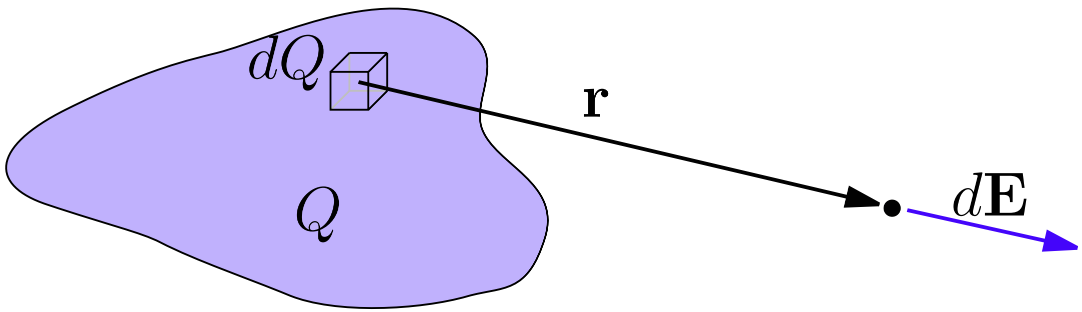

  - If small enough, each chunk obeys Coulomb's law

$$d\mathbf{E} = \frac{k\,dQ}{r^2}\hat{\mathbf{r}}.$$

Continuous charge distributions, contd

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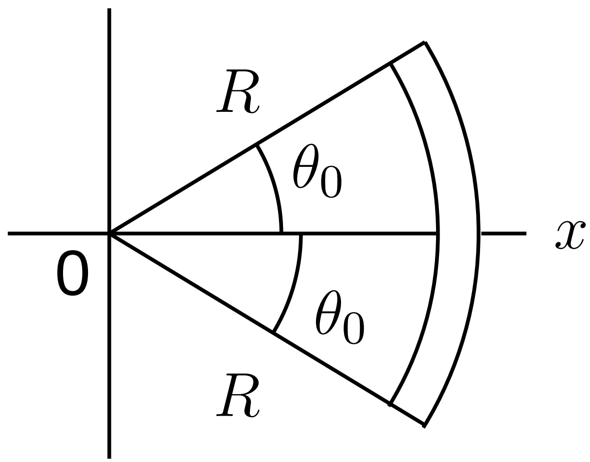
  - Method can be difficult but is guaranteed to work. Next class will show easier method that works in special cases.

Continuous charge distributions, contd

- **Example: Pr. 49**
  - A thin rod bent into the shape of an arc of a circle of radius R carries a uniform charge per unit length λ. The arc subtends a total angle $2\theta_0$, symmetric about the x axis, as shown below. Determine the electric field $\mathbf{E}$ at the origin 0.



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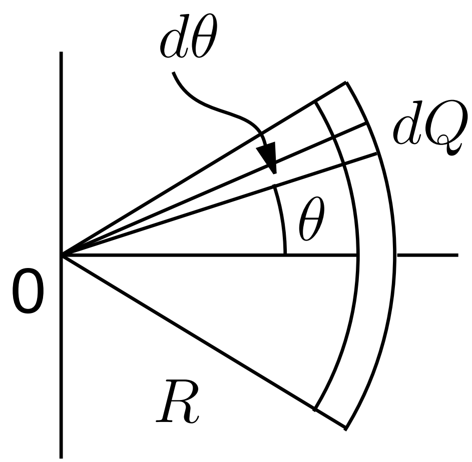

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  - We need to divide this continuous charge distribution into discrete chunks of charge dQ.
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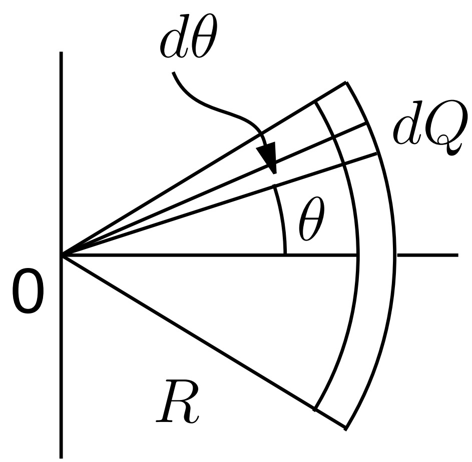

  - Then the chunk has charge $dQ = \lambda R\, d\theta$.

Continuous charge distributions, contd

- **Solution: Pr. 49, contd**

Continuous charge distributions, contd

- **Solution: Pr. 49, contd**
  - Now we can apply Coulomb's law to get the electric field due to a single chunk,

$$d\mathbf{E} = \frac{k\,dQ}{R^2}(-\cos\theta\,\hat{\mathbf{i}} - \sin\theta\,\hat{\mathbf{j}}).$$

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$$d\mathbf{E} = \frac{k\,dQ}{R^2}(-\cos\theta\,\hat{\mathbf{i}} - \sin\theta\,\hat{\mathbf{j}}).$$

  - We can sum over all the chunks to get the total electric field,

$$\begin{aligned}\mathbf{E} &= \int d\mathbf{E} \\ &= \frac{k}{R^2}\int dQ(-\cos\theta\,\hat{\mathbf{i}} - \sin\theta\,\hat{\mathbf{j}})\end{aligned}$$

Continuous charge distributions, contd

- **Solution: Pr. 49, contd**

$$\begin{aligned} \mathbf{E} &= \frac{k}{R^2}\int_{-\theta_0}^{+\theta_0} \lambda R\, d\theta(-\cos\theta\,\hat{\mathbf{i}} - \sin\theta\,\hat{\mathbf{j}}) \\ &= -\frac{k\lambda}{R}\left[\hat{\mathbf{i}}\int_{-\theta_0}^{+\theta_0}\cos\theta\, d\theta + \hat{\mathbf{j}}\int_{-\theta_0}^{+\theta_0}\sin\theta\, d\theta\right]. \end{aligned}$$

Continuous charge distributions, contd

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$$
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\mathbf{E} &= \frac{k}{R^2}\int_{-\theta_0}^{+\theta_0} \lambda R\, d\theta(-\cos\theta\,\hat{\mathbf{i}} - \sin\theta\,\hat{\mathbf{j}}) \\
&= -\frac{k\lambda}{R}\left[\hat{\mathbf{i}}\int_{-\theta_0}^{+\theta_0}\cos\theta\, d\theta + \hat{\mathbf{j}}\int_{-\theta_0}^{+\theta_0}\sin\theta\, d\theta\right].
\end{aligned}
$$

- Notice the $\hat{\mathbf{j}}$ contributions cancel out, leaving only

$$
\begin{aligned}
\mathbf{E} &= -\frac{k\lambda}{R}\,\hat{\mathbf{i}}\,2\sin\theta_0 \\
&= -\frac{2k\lambda\sin\theta_0}{R}\hat{\mathbf{i}}. \quad \square
\end{aligned}
$$

End

- **Practice Problems:**
  - Ch. 21: Q. 15, 17, 19, 21, 23
  - Ch. 21: Pr. 11, 13, 15, 19, 25, 27, 29, 35, 37, 39, 41, 43, 55, 57, 71, 73, 75, 77, 79, 81, 83, 87

End

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- **Interactive Quiz: Feedback**

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- **Interactive Quiz: Feedback**
- **Tutorial Question: tut05**