533D: Animation Physics

- http://www.cs.ubc.ca/~rbridson/courses/ 533d-winter-2005
- Course schedule

Slides online, but you need to take notes too!

- Reading
 - Relevant animation papers as we go
- Assignments + Final Project information
- Resources

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Contacting Me

- Robert Bridson
 - CICSR 189 (moving to CS2 in reading week)
 - Drop by, or make an appointment
 - 604-822-1993 (or just 21993)
 - email rbridson@cs.ubc.ca
 - Newsgroup ubc.courses.cpsc.533b
- I always like feedback!

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Evaluation

- 4 assignments (60%)
 - See the web for details + when they are due
 - Mostly programming, with a little analysis (writing)
- Also a final project (40%)
 - Details will come later, but basically you need to either significantly extend an assignment or animate something else - talk to me about topics
 - Present in final class informal talk, show movies
- Late: without a good reason, 20% off per day
- For final project starts after final class
 - For assignments starts morning after due

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Why?

- Natural phenomena: passive motion
- Film/TV: difficult with traditional techniques
 When you control every detail of the motion, it's had
 - When you control every detail of the motion, it's hard to make it look like it's not being controlled!
- Games: difficult to handle everything convincingly with prescripted motion
- Computer power is increasing, audience expectations are increasing, artist power isn't: need more automatic methods
- Directly simulate the underlying physics to get realistic motion

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Topics

- Particle Systems
 - the basics
- Deformable Bodies
 - e.g. cloth and flesh
- Constrained Dynamics
 - e.g. rigid bodies
- ♦ Fluids
 - e.g. water

Particle Systems

- ♦ Read:
- Reeves, "Particle Systems...", SIGGRAPH'83 Sims, "Particle animation and rendering using data parallel computation", SIGGRAPH '90
- Some phenomena is most naturally described as many small particles
 - Rain, snow, dust, sparks, gravel, ...
- Others are difficult to get a handle on
 - Fire, water, grass, ...

Particle Basics

- Each particle has a position
 - Maybe orientation, age, colour, velocity, temperature, radius, ...
 - Call the state x
- Seeded randomly somewhere at start
 - Maybe some created each frame
- Move (evolve state x) each frame according to some formula
- Eventually die when some condition met

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Example

- Sparks from a campfire
- Every frame (1/24 s) add 2-3 particles
 - Position randomly in fire
 - Initialize temperature randomly
- Move in specified turbulent smoke flow
 Also decrease temperature
- Render as a glowing dot (blackbody radiation from temperature)
- Kill when too cold to glow visibly

Rendering

- We won't talk much about rendering in this course, but most important for particles
- The real strength of the idea of particle systems: how to render
 - Could just be coloured dots
 - Or could be shards of glass, or animated sprites (e.g. fire), or deforming blobs of water, or blades of grass, or birds in flight, or ...

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First Order Motion

First Order Motion

 For each particle, have a simple 1st order differential equation:

$$\frac{dx}{dt} = v(x,t)$$

- Analytic solutions hopeless
- Need to solve this numerically forward in time from x(t=0) to
 - x(frame1), x(frame2), x(frame3), ...
 - May be convenient to solve at some intermediate times between frames too

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Forward Euler

Simplest method:

Or:

$$\frac{x_{n+1} - x_n}{\Delta t} = v(x_n, t_n)$$
$$x_{n+1} = x_n + \Delta t v(x_n, t_n)$$

- Can show it's first order accurate:
 Error accumulated by a fixed time is O(Δt)
- Thus it converges to the right answer
 - Do we care?

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Aside on Error

- General idea want error to be small
 - Obvious approach: make Δt small
 - But then need more time steps expensive
- Also note O(1) error made in modeling
 - Even if numerical error was 0, still wrong!
 - In science, need to validate against experiments
 - In graphics, the experiment is showing it to an audience: does it look real?
- So numerical error can be huge, as long as your solution has the right qualitative look

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Forward Euler Stability

- Big problem with Forward Euler: it's not very stable
- Example: dx/dt = -x, x(0) = 1
- ♦ Real solution e^{-t} smoothly decays to zero, always positive
- Run Forward Euler with Δt=11
 - x=1, -10, 100, -1000, 10000, ...
 - Instead of 1, 1.7*10⁻⁵, 2.8*10⁻¹⁰, ...

Linear Analysis

Approximate

$$v(x,t) \approx v(x^*,t^*) + \frac{\partial v}{\partial x} \cdot (x-x^*) + \frac{\partial v}{\partial t} \cdot (t-t^*)$$

 Ignore all but the middle term (the one that could cause blow-up)

dx/dt = Ax

• Look at x parallel to eigenvector of A: the "test equation" $dx/dt = \lambda x$

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The Test Equation

- Get a rough, hazy, heuristic picture of the stability of a method
- Note that eigenvalue λ can be complex
- But, assume that for real physics
 - Things don't blow up without bound
 - Thus real part of eigenvalue λ is ≤ 0
- Beware!
 - · Nonlinear effects can cause instability
 - Even with linear problems, what follows assumes constant time steps - varying (but supposedly stable) steps can induce instability
 - see J. P. Wright, "Numerical instability due to varying time steps...", JCP 1998

Using the Test Equation

Forward Euler on test equation is

$$x_{n+1} = x_n + \Delta t \,\lambda x_n$$

Solving gives

$$x_n = \left(1 + \lambda \Delta t\right)^n x_0$$

◆ So for stability, need

 $|1 + \lambda \Delta t| < 1$

• Can plot all the values of $\lambda\Delta t$ on the complex plane where F.E. is stable:



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Real Eigenvalue

- Say eigenvalue is real (and negative)
 - Corresponds to a damping motion, smoothly coming to a halt 2

$$\Delta t < \frac{2}{\lambda}$$

- Is this bad?
 - If eigenvalue is big, could mean small time steps
 - · But, maybe we really need to capture that time scale anyways, so no big deal

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Imaginary Eigenvalue

- If eigenvalue is pure imaginary...
 - Oscillatory or rotational motion
- Cannot make ∆t small enough
- Forward Euler unconditionally unstable for these kinds of problems!
- Need to look at other methods

Runge-Kutta Methods

- Also "explicit"
 - next x is an explicit function of previous
- But evaluate v at a few locations to get a better estimate of next x
- E.g. midpoint method (one of RK2)

$$x_{n+\frac{1}{2}} = x_n + \frac{1}{2}\Delta t v(x_n, t_n)$$

$$x_{n+1} = x_n + \Delta t v(x_{n+\frac{1}{2}}, t_{n+\frac{1}{2}})$$

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Midpoint RK2

- Second order: error is O(Δt²) when smooth
- Larger stability region:



• But still not stable on imaginary axis: no point

Modified Euler

- (Not an official name)
- Lose second-order accuracy, get stability on imaginary axis:

$$x_{n+\alpha} = x_n + \alpha \Delta t v(x_n, t_n)$$
$$x_{n+1} = x_n + \Delta t v(x_{n+\alpha}, t_{n+\alpha})$$

• Parameter α between 0.5 and 1 gives trade-off between imaginary axis and real axis

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Modified Euler (2)

• Stability region for α =2/3



- υ Great! But twice the cost of Forward Euler
- Can you get more stability per vevaluation?

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Higher Order Runge-Kutta

 RK3 and up naturally include part of the imaginary axis



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TVD-RK3

 RK3 useful because it can be written as a combination of Forward Euler steps and averaging: can guarantee stuff!

$$\begin{split} \tilde{x}_{n+1} &= x_n + \Delta t v (x_n, t_n) \\ \tilde{x}_{n+2} &= \tilde{x}_{n+1} + \Delta t v (\tilde{x}_{n+1}, t_{n+1}) \\ \tilde{x}_{n+\frac{1}{2}} &= \frac{3}{4} x_n + \frac{1}{4} \tilde{x}_{n+2} \\ \tilde{x}_{n+\frac{3}{2}} &= \tilde{x}_{n+\frac{1}{2}} + \Delta t v (\tilde{x}_{n+\frac{1}{2}}, t_{n+\frac{1}{2}}) \\ x_{n+1} &= \frac{1}{3} x_n + \frac{2}{3} \tilde{x}_{n+\frac{3}{2}} \end{split}$$

RK4

Often most bang for the buck

 $\begin{aligned} v_1 &= v \Big(x_n, t_n \Big) \\ v_2 &= v \Big(x_n + \frac{1}{2} \Delta t v_1, t_{n+\frac{1}{2}} \Big) \\ v_3 &= v \Big(x_n + \frac{1}{2} \Delta t v_2, t_{n+\frac{1}{2}} \Big) \\ v_4 &= v \Big(x_n + \Delta t v_3, t_{n+1} \Big) \\ x_{n+1} &= x_n + \Delta t \Big(\frac{1}{6} v_1 + \frac{2}{6} v_2 + \frac{2}{6} v_3 + \frac{1}{6} v_4 \Big) \end{aligned}$

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Selecting Time Steps

Selecting Time Steps

- Hack: try until it looks like it works
- Stability based:
 - · Figure out a bound on magnitude of Jacobian
 - Scale back by a fudge factor (e.g. 0.9, 0.5)
 Try until it looks like it works... (remember all the dubious assumptions we made for linear stability analysis!)
 - Why is this better than just hacking around in the first place?
- Adaptive error based:
 - Usually not worth the trouble in graphics

Time Stepping

- Sometimes can pick constant ∆t
 - One frame, or 1/8th of a frame, or ...
- Often need to allow for variable ∆t
 - Changing stability limit due to changing Jacobian
 - Difficulty in Newton converging
 - ...
- But prefer to land at the exact frame time
 - So clamp Δt so you can't overshoot the frame

Example Time Stepping Algorithm

- Set done = false
- While not done
 - Find good Δt
 - If $t+\Delta t \ge t_{frame}$
 - Set ∆t = t_{frame}-t
 - Set done = true
 - Else if t+1.5∆t ≥ t_{frame}
 - Set ∆t = 0.5(t_{frame}-t)
 - ...process time step...
 - Set t = t+∆t
- Write out frame data, continue to next frame

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Implicit Methods

Large Time Steps

- Look at the test equation $\frac{dx}{dt} = \lambda x$
- Exact solution is $x(t_{n+1}) = e^{\lambda \Delta t} x(t_n) = \left(1 + \lambda \Delta t + \frac{1}{2} (\lambda \Delta t)^2 + \ldots\right) x(t_n)$
- Explicit methods approximate this with polynomials (e.g. Taylor)
- Polynomials must blow up as t gets big
 Hence explicit methods have stability limit
- We may want a different kind of approximation that drops to zero as ∆t gets big
 - Avoid having a small stability limit when error says it should be fine to take large steps ("stiffness")

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Simplest stable approximation

- Instead use $e^{\lambda \Delta t} \approx \frac{1}{1 \lambda \Delta t}$
- That is, $x_{n+1} = \frac{1}{1 \lambda \Delta t} x_n$
- Rewriting: $x_{n+1} = x_n + \Delta t \lambda x_{n+1}$
- This is an "implicit" method: the next x is an implicit function of the previous x
 - Need to solve equations to figure it out

Backward Euler

• The simplest implicit method:

$$x_{n+1} = x_n + \Delta t v (x_{n+1}, t_{n+1})$$

- First order accurate
- Test equation shows stable when $|1 \lambda \Delta t| > 1$
- This includes everything except a circle in the positive real-part half-plane
- It's stable even when the physics is unstable!
- This is the biggest problem: damps out motion unrealistically

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Aside: Solving Systems

- If v is linear in x, just a system of linear equations
 - If very small, use determinant formula
 - If small, use LAPACK
 - If large, life gets more interesting...
- If v is mildly nonlinear, can approximate with linear equations ("semi-implicit")

$$\begin{aligned} x_{n+1} &= x_n + \Delta t \, v(x_{n+1}) \\ &\approx x_n + \Delta t \bigg(v(x_n) + \frac{\partial v(x_n)}{\partial x} (x_{n+1} - x_n) \bigg) \end{aligned}$$

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Newton's Method

- For more strongly nonlinear v, need to iterate:
 - Start with guess x_n for x_{n+1} (for example)
 - Linearize around current guess, solve linear system for next guess
 - Repeat, until close enough to solved
- Note: Newton's method is great when it works, but it might not work
 - If it doesn't, can reduce time step size to make equations easier to solve, and try again
 - Maybe use higher power optimization methods (e.g. at least use line search)

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Newton's Method: B.E.

- Start with x⁰=x_n (simplest guess for x_{n+1})
- For k=1, 2, ... find $x^{k+1}=x^k+\Delta x$ by solving

$$x^{k+1} = x_n + \Delta t \left(v(x^k) + \frac{\partial v(x^k)}{\partial x} (x^{k+1} - x^k) \right)$$
$$\Rightarrow \left(I - \Delta t \frac{\partial v(x^k)}{\partial x} \right) \Delta x = x_n + \Delta t v(x^k) - x^k$$

- To include line-search for more robustness, change update to x^{k+1}=x^k+αΔx and choose 0 < α ≤ 1 that minimizes ||x_n + Δtν(x^{k+1}, t_{n+1}) x^{k+1}||
- Stop when right-hand side is small enough, set x_{n+1}=x^k

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Trapezoidal Rule

Can improve by going to second order:

$$x_{n+1} = x_n + \Delta t \left(\frac{1}{2} v(x_n, t_n) + \frac{1}{2} v(x_{n+1}, t_{n+1}) \right)$$

- This is actually just a half step of F.E., followed by a half step of B.E.
 - F.E. is under-stable, B.E. is over-stable, the combination is just right
- Stability region is the left half of the plane: exactly the same as the physics!
- Really good for pure rotation (doesn't amplify or damp)

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