

Notes

- ◆ Assignment questions...

Linear Triangle Elements

- ◆ Let's work out the formulas...

The equations

$$\int_{\Omega} \phi_j \sum_i \rho \ddot{x}_i \phi_i - \int_{\Omega} \phi_j f_{body} + \int_{\Omega} \sigma \nabla \phi_j = 0$$

$$\sum_i \int_{\Omega} \rho \phi_j \phi_i \ddot{x}_i = \int_{\Omega} \phi_j f_{body} - \int_{\Omega} \sigma \nabla \phi_j$$

- Note that ϕ_j is zero on all but the triangles surrounding j , so integrals simplify
- Also: naturally split integration into separate triangles

Change in momentum term

- ◆ Let $m_{ij} = \int \rho \phi_i \phi_j$
- ◆ Then the first term is just $\sum_i m_{ji} \ddot{x}_i$
- ◆ Let $M=[m_{ij}]$: then first term is $M \ddot{x}$
- ◆ M is called the mass matrix
 - Obviously symmetric (actually SPD)
 - Not diagonal!
- ◆ Note that once we have the forces (the other integrals), we need to invert M to get accelerations

Body force term

- ◆ Usually just gravity: $f_{body} = \rho g$
- ∪ Rather than do the integral with density all over again, use the fact that ϕ_i sum to 1
 - They form a "partition of unity"
 - They represent constant functions exactly - just about necessary for convergence
- ◆ Then body force term is $gM1$
- ◆ More specifically, can just add g to the accelerations; don't bother with integrals or mass matrix at all

Stress term

- ◆ Calculate constant strain and strain rate (so constant stress) for each triangle separately
- ◆ Note $\nabla \phi_j$ is constant too
- ∪ So just take $\sigma \nabla \phi_j$ times triangle area
- ∪ [derive what $\nabla \phi_j$ is]
- ∪ Magic: exact same as FVM!
 - In fact, proof of convergence of FVM is often (in other settings too) proved by showing it's equivalent or close to some kind of FEM

The algorithm

- ◆ Loop over triangles
 - Loop over corners
 - Compute integral terms
 - only need to compute M once though - it's constant
 - End up with row of M and a “force”
- ◆ Solve $Ma=f$
- ◆ Plug this a into time integration scheme

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Lumped Mass

- ◆ Inverting mass matrix unsatisfactory
 - For particles and FVM, each particle had a mass, so we just did a division
 - Here mass is spread out, need to do a big linear solve - even for explicit time stepping
- ◆ Idea of lumping: replace M with the “lumped mass matrix”
 - A diagonal matrix with the same row sums
 - Inverting diagonal matrix is just divisions - so diagonal entries of lumped mass matrix are the particle masses
 - Equivalent to FVM with centroid-based volumes

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Consistent vs. Lumped

- ◆ Original mass matrix called “consistent”
- ◆ Turns out its strongly diagonal dominant (fairly easy to solve)
- ◆ Multiplying by mass matrix = smoothing
- ◆ Inverting mass matrix = sharpening
- ◆ Rule of thumb:
 - Implicit time stepping - use consistent M (counteract over-smoothing, solving system anyways)
 - Explicit time stepping - use lumped M (avoid solving systems, don't need extra sharpening)

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Locking

- ◆ Simple linear basis actually has a major problem: locking
 - But graphics people still use them all the time...
- ◆ Notion of numerical stiffness
 - Instead of thinking of numerical method as just getting an approximate solution to a real problem,
 - Think of numerical method as exactly solving a problem that's nearby
 - For elasticity, we're exactly solving the equations for a material with slightly different (and not quite homogeneous/isotropic) stiffness
- ◆ Locking comes up when numerical stiffness is MUCH higher than real stiffness

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Locking and linear elements

- ◆ Look at nearly incompressible materials
- ◆ Can a linear triangle mesh deform incompressibly?
 - [derive problem]
- ◆ Then linear elements will resist far too much: numerical stiffness much too high
- ◆ Numerical material “locks”
- ◆ FEM isn't really a black box!
- ◆ Solutions:
 - Don't do incompressibility
 - Use other sorts of elements (quads, higher order)

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Quadrature

- ◆ Formulas for linear triangle elements and constant density simple to work out
- ◆ Formulas for subdivision surfaces (or high-order polynomials, or splines, or wavelets...) and varying density are NASTY
- ◆ Instead use “quadrature”
 - I.e. numerical approximation to integrals
- ◆ Generalizations of midpoint rule
 - E.g. Gaussian quadrature (for intervals, triangles, tets) or tensor products (for quads, hexes)
- ◆ Make sure to match order of accuracy [or not]

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Accuracy

- ◆ At least for SPD linear problems (e.g. linear elasticity) FEM selects function from finite space that is “closest” to solution
 - Measured in a least-squares, energy-norm sense
- ◆ Thus it’s all about how well you can approximate functions with the finite space you chose
 - Linear or bilinear elements: $O(h^2)$
 - Higher order polynomials, splines, etc.: better

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Hyper-elasticity

- ◆ Another common way to look at elasticity
 - Useful for handling weird nonlinear compressibility laws, for reduced dimension models, and more
- ◆ Instead of defining stress, define an elastic potential energy
 - Strain energy density $W=W(A)$
 - $W=0$ for no deformation, $W>0$ for deformation
 - Total potential energy is integral of W over object
- ◆ This is called hyper-elasticity or Green elasticity
- ◆ For most (the ones that make sense) stress-strain relationships can define W
 - E.g. linear relationship: $W=\sigma:\varepsilon=\text{trace}(\sigma^T\varepsilon)$

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Variational Derivatives

- ◆ Force is the negative gradient of potential
 - Just like gravity
- ◆ What does this mean for a continuum?
 - $W=W(\partial X/\partial p)$, how do you do $-d/dX$?
- ◆ Variational derivative: $W_{total}[X + \varepsilon Y] = \int W\left(\frac{\partial X}{\partial p} + \varepsilon \frac{\partial Y}{\partial p}\right)$
 - So variational derivative is $-\nabla \cdot \partial W / \partial A$
 - And $f = \nabla \cdot \partial W / \partial A$
 - Then stress is $\partial W / \partial A$

$$\begin{aligned} &\approx \int W\left(\frac{\partial X}{\partial p}\right) + \varepsilon \frac{\partial W}{\partial A} \frac{\partial Y}{\partial p} \\ &= W_{total} + \varepsilon \int \frac{\partial W}{\partial A} \frac{\partial Y}{\partial p} \\ &= W_{total} - \varepsilon \int Y \nabla \cdot \frac{\partial W}{\partial A} \end{aligned}$$

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Numerics

- ◆ Simpler approach: find discrete W_{total} as a sum of W 's for each element
 - Evaluate just like FEM, or any way you want
- ◆ Take gradient w.r.t. positions $\{x_i\}$
 - Ends up being a Galerkin method
- ◆ Also note that an implicit method might need Jacobian = negative Hessian of energy
 - Must be symmetric, and at least near stable configurations must be negative definite

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Curve / Springs

- ◆ Take $W(A)=1/2 E(|A|-1)^2 L$ for each segment
 - Note factor of L : this is approximation to an integral over segment in object space of length L
- ◆ $A=(x_{i+1}-x_i)/L$ is the deformation gradient for piecewise linear elements
- ◆ Then take derivative w.r.t. x_i to get this element’s contribution to force on i
- ◆ Lo and behold [exercise] get exactly the original spring force
- ◆ Note: defining stress and strain would be more complicated, because of the dimension differences
 - A is 3×1 , not square

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