

Notes

- Please read:
 - Choi & Ko, “Stable but responsive cloth”, SIGGRAPH’02
 - Grinspun et al., “Discrete shells”, SCA’03
- Homework #4 (cloth simulation) goes out today
 - Due March 4, but you may want to look at it over the break

Interlude...

- Linearize and simplify drastically, look for steady-state solution ($F=0$): spline equations
 - Essentially 4th derivatives are zero
 - Solutions are (bi-)cubics
- Model (nonsteady) problem: $x_{tt} = -x_{pppp}$
 - Assume solution $x(p, t) = e^{\sqrt{-1}k(p-ct)}$

Wave of spatial frequency k , moving at speed c

 - [solve for wave parameters]
 - Dispersion relation: small waves move really fast
 - CFL limit (and stability): for fine grids, BAD
 - Thankfully, we rarely get that fine

Bending energy

- Bending is very difficult to get a handle on without variational approach
- Bending strain energy density:
 $W = 1/2 B \kappa^2$
- Here κ is mean curvature
 - Look at circles that fit surface
 - Maximum radius R and minimum radius r
 - $\kappa = (1/R + 1/r)/2$
 - Can define directly from second derivatives of $X(p)$
 - Uh-oh - second derivatives? [FEM nastier]
 - W is 2nd order, stress is 3rd order, force is 4th order derivatives!

Implicit/Explicit Methods

- Implicit bending is painful
- In graphics, usually unnecessary
 - Dominant forces on the grid resolution we use tend to be the 2nd order terms: stretching etc.
- But nice to go implicit to avoid time step restriction for stretching terms
- No problem: treat some terms (bending) explicitly, others (stretching) implicitly
 - $v_{n+1} = v_n + \Delta t / m (F_1(x_n, v_n) + F_2(x_{n+1}, v_{n+1}))$
 - All bending is in F_1 , half the elastic stretch in F_1 , half the elastic stretch in F_2 , all the damping in F_2

Discrete Mean Curvature

- [draw triangle pair]
- κ for that chunk varies as $\kappa \sim \frac{\theta}{h_1 + h_2}$
- So integral of κ^2 varies as $W = \sum_e \frac{\theta^2}{(h_1 + h_2)^2} (|\Delta_1| + |\Delta_2|)$
 $\sim \sum_e \frac{\theta^2 |e|^2}{|\Delta_1| + |\Delta_2|}$
- Edge length, triangle areas, normals are all easy to calculate
- θ needs inverse trig functions
- But θ^2 behaves a lot like $1 - \cos(\theta/2)$ over interval $[-\pi, \pi]$ [draw picture]

Gradient of Theta

- Can use implicit differentiation on $\cos(\theta) = n_1 \cdot n_2$
 - Not too fun
- Another approach: modal analysis
 - Gradient is orthogonal to isocontours of theta
 - Find a basis for tangent plane to isocontours (perpendicular to grad(theta))
 - [go through 11 modes: 6 rigid body modes, 5 planar deformations]
 - Solve for the mode that is orthogonal to all the rest
 - This gives directions, do a little diff. for mag

Bending Force

- Force on x_i due to bending element involving i is then $F_i = -B \frac{\partial W}{\partial x_i} \sim -B \frac{|e|^2}{|\Delta_1| + |\Delta_2|} \sin\left(\frac{\theta}{2}\right) \frac{\partial \theta}{\partial x_i}$
- Treat first terms as a constant (precompute in the rest configuration)

$$\sin \frac{\theta}{2} = \pm \sqrt{\frac{1}{2}(1 - n_1 \cdot n_2)}$$
- Sign should be the same as $\sin \theta = n_1 \times n_2 \cdot \hat{e}$
- Still need to compute $\partial \theta / \partial x_i$

The bending mode

- Vertices 1 to 4 with common edge $e = x_4 - x_3$
 - [see Bridson et al., "Simulation of clothing...", SCA'03]
 - $u_i = \partial \theta / \partial x_i$ $\Delta_1 = (x_1 - x_3) \times (x_1 - x_4)$ $\Delta_2 = (x_2 - x_4) \times (x_2 - x_3)$

$$u_1 = |e| \frac{\Delta_1}{|\Delta_1|^2} \quad u_2 = |e| \frac{\Delta_2}{|\Delta_2|^2}$$

$$u_3 = \frac{(x_1 - x_4) \cdot e}{|e|} \frac{\Delta_1}{|\Delta_1|^2} + \frac{(x_2 - x_4) \cdot e}{|e|} \frac{\Delta_2}{|\Delta_2|^2}$$

$$u_4 = -\frac{(x_1 - x_3) \cdot e}{|e|} \frac{\Delta_1}{|\Delta_1|^2} - \frac{(x_2 - x_3) \cdot e}{|e|} \frac{\Delta_2}{|\Delta_2|^2}$$

Modal Analysis and Soft Constraints

- Go back to $W = \frac{1}{2}kC(x)^T C(x)$
- Then elastic force is $F = -\frac{\partial W}{\partial x} = -k \frac{\partial C^T}{\partial x} C$
- Here $C=\theta$
- Think of $\partial C/\partial x$ as the mode which directly influences C (steepest ascent)
 - Any other direction moves along isocontours of C: inefficient, interferes with other physics
- So we want force in this direction
- And the magnitude should be proportional to $-C$ to restore C to 0

Damping bending

- Follow the same reasoning:

$$F = -D|e|uu^T v = -D|e|u\dot{\theta}$$
- The lcl factor is to get it to converge as the mesh is refined

Damping force

- We also want damping force to be in the same mode
 - Any other direction would again interfere with other physics
- And should be proportional to the component of the velocities in this mode
 - The other components of velocity shouldn't influence the damping
- So damping force is

$$F = -D \frac{\partial C^T}{\partial x} \frac{\partial C}{\partial x} v = -D \frac{\partial C^T}{\partial x} \dot{C}$$

- Note symmetric, negative semi-definite, linear in v

Shells

- “Plates” in elasticity refer to surfaces that resist bending, and have a flat rest pose
- “Shells” are same but with a curved rest pose
- We can easily do shells by storing a nonzero rest angle
- Replace $\sin(\theta/2)$ in elastic force by $\sin(\theta/2) - \sin(\theta_0/2)$
 - Or if θ_0 is far from zero, may need to go to $\theta - \theta_0$
- Note: using second neighbour springs you can't do shells! [popping - also with inconsistent edge springs]

Other stuff in cloth mechanics

- Better model planar elasticity
 - Take into account anisotropy of weave and weft
- Creasing - allow rest angle to move when current angle too big
 - Plasticity coming up after the break...
- Subdivision for rendering
 - Geometric buckling, collision-aware
- Adaptive meshes
- Interaction with the wind
- Speed-ups

Not quite...

- For large collision steps this can be bad
 - This is not just a particle system: particles are connected!
 - If a collision impulse discontinuously bounces one particle off, nearby cloth should feel it
 - Otherwise can have excessive strains (or strain rates) to deal with next time step
 - Want to naturally smooth out discontinuity in velocity field

Cloth collisions

- Can use just inelastic collisions
 - Continuum takes care of rebound naturally
- Cloth colliding against non-sharp thick objects
 - Just treat as a particle system, like assignment 2
 - Can even allow particles to interpenetrate a bit (staying relaxed about this can let cloth look nicer)
 - For rendering need to push cloth -- the triangles as well as the vertices -- out of the objects
 - [wrinkle-preserving map]
- Cloth colliding against non-sharp thin objects
 - Need to worry about robustness -- solve for collision times, use repulsions.

Implicit Velocity Smoothing

- If we just evolve $v_t = F_{\text{damp}}/m$ we get smoothing
 - Damping forces seek to minimize relative velocities
 - Model problem is $v_t = v_{xx}$, the heat equation (also in multiple dimensions)
 - Solution is convolution with a Gaussian
 - Doing one implicit time step of this equation is a similar smoothing
 - Explicit time steps can only do very local smoothing...
- We'll use this to filter velocities after collision processing

Collision time steps for cloth

- Evolve (x,v) to candidate new state
- Check for collisions
 - Use average velocity $v^{avg}=(x^{new}-x)/\Delta t$ for collision formulas (only doing inelastic)
 - Resolve with a collision resolution pipeline
- Get x^{new} with collisions resolved
- Find updated average velocity $v^{avg}=(x^{new}-x)/\Delta t$
 - This “lives” at the midpoint: $v(t+\Delta t/2)$
- Use implicit velocity update to get end value:
 - Note it’s just linear
$$v_{n+1} = v_{n+1/2} + \frac{\Delta t}{2} M^{-1} F(x_{n+1}, v_{n+1})$$

Self-collision

- Need a collision resolution pipeline like we had for rigid bodies
 - Simple, accurate, fast but local algorithms at the start
 - Robust, global algorithms at the end
- Robustness is useful: it’s not so easy to recover from a mistake without harming motion
 - Can do it though, and might need to (animators doing the impossible...)
 - See Baraff et al, SIGGRAPH’03