#### **Notes**

- Most of assignment 1 hasn't been covered in class yet, but after today you should be able to do a lot of it
- Forgot to include instructions about view\_obj:
  - To navigate, hold down shift and click/drag with left, right, or middle mouse buttons (same navigation model as Maya)

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# **Trapezoidal Rule Again**

◆ The method:

$$x_{n+1} = x_n + \Delta t \left( \frac{1}{2} v(x_n, t_n) + \frac{1}{2} v(x_{n+1}, t_{n+1}) \right)$$

◆ Let's work out stability:

$$x_{n+1} = x_n + \Delta t \left(\frac{1}{2} \lambda x_n + \frac{1}{2} \lambda x_{n+1}\right)$$
$$\left(1 - \frac{1}{2} \lambda \Delta t\right) x_{n+1} = \left(1 + \frac{1}{2} \lambda \Delta t\right) x_n$$
$$x_{n+1} = \frac{1 + \frac{1}{2} \lambda \Delta}{1 - \frac{1}{2} \lambda \Delta t} x_n$$

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# **Monotonicity**

- Test equation with real, negative  $\lambda$ 
  - True solution is x(t)=x<sub>0</sub>e<sup>λt</sup>, which smoothly decays to zero, doesn't change sign (monotone)
- ◆ Forward Euler at stability limit:
  - x=x<sub>0</sub>, -x<sub>0</sub>, x<sub>0</sub>, -x<sub>0</sub>, ...
- ◆ Not smooth, oscillating sign: garbage!
- So monotonicity limit stricter than stability in this case
- ◆ RK3 has the same problem
  - But the even order RK are fine for linear problems
  - TVD-RK3 designed so that it's fine when F.E. is, even for nonlinear problems!

# **Monotonicity and Implicit Methods**

- Backward Euler is unconditionally monotone
  - No problems with oscillation, just too much damping
- ◆ Trapezoidal Rule suffers though, because of that half-step of F.E.
  - Beware: could get ugly oscillation instead of smooth damping

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### **Summary 1**

- ◆ Particle Systems: useful for lots of stuff
- ◆ Need to move particles in velocity field
- ◆ Forward Euler
  - Simple, first choice unless problem has oscillation/rotation
- ◆ Runge-Kutta if happy to obey stability limit
  - Modified Euler may be cheapest method
  - RK4 general purpose workhorse
  - TVD-RK3 for more robustness with nonlinearity (more on this later in the course!)

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### **Summary 2**

- ◆ If stability limit is a problem, look at implicit methods
  - e.g. need to guarantee a frame-rate, or explicit time steps are way too small
- ◆ Trapezoidal Rule
  - If monotonicity isn't a problem
- ◆ Backward Euler
  - Almost always works, but may over-damp!

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#### **Second Order Motion**

#### **Second Order Motion**

- ◆ If particle state is just position (and colour, size, ...) then 1st order motion
  - No inertia
  - Good for very light particles that stay suspended: smoke, dust...
  - Good for some special cases (hacks)
- But most often, want inertia
  - State includes velocity, specify acceleration
  - Can then do parabolic arcs due to gravity, etc.
- This puts us in the realm of standard Newtonian physics
  - F=ma
- Alternatively put:
  - dx/dt=v
  - dv/dt=F(x,v,t)/m (i.e. a(x,v,t))
- ◆ For systems (with many masses) say dv/dt=M<sup>-1</sup>F(x,v,t) where M is the "mass matrix" masses on the diagonal

#### What's New?

- ♦ If x=(x,v) this is just a special form of 1st order: dx/dt=v(x,t)
- But since we know the special structure, can we take advantage of it? (i.e. better time integration algorithms)
  - More stability for less cost?
  - Handle position and velocity differently to better control error?

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# **Linear Analysis**

Approximate acceleration:

$$a(x,v) \approx a_0 + \frac{\partial a}{\partial x}x + \frac{\partial a}{\partial v}v$$

- ◆ Split up analysis into different cases
- ◆ Begin with first term dominating: constant acceleration
  - e.g. gravity is most important

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## **Constant Acceleration**

- ♦ Solution is  $v(t) = v_0 + a_0 t$  $x(t) = x_0 + v_0 t + \frac{1}{2} a_0 t^2$
- No problem to get v(t) right: just need 1st order accuracy
- ◆ But x(t) demands 2nd order accuracy
- ◆ So we can look at mixed methods:
  - 1st order in v
  - 2nd order in x

#### **Linear Acceleration**

- ◆ Dependence on x and v dominates: a(x,v)=-Kx-Dv
- ◆ Do the analysis as before:

$$\frac{d}{dt} \begin{pmatrix} x \\ v \end{pmatrix} = \begin{pmatrix} 0 & I \\ -K & -D \end{pmatrix} \begin{pmatrix} x \\ v \end{pmatrix}$$

◆ Eigenvalues of this matrix?

# **More Approximations...**

- Typically K and D are symmetric semi-definite (there are good reasons)
  - What does this mean about their eigenvalues?
- Often, D is a linear combination of K and I ("Rayleigh damping"), or at least close to it
  - Then K and D have the same eigenvectors (but different eigenvalues)
  - Then the eigenvectors of the Jacobian are of the form  $(u, \alpha u)^T$
  - [work out what  $\alpha$  is in terms of  $\lambda_K$  and  $\lambda_D$ ]

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# **Simplification**

 $\nu$   $\alpha$  is the eigenvalue of the Jacobian, and

$$\alpha = -\frac{1}{2} \lambda_D \pm \sqrt{\left(\frac{1}{2} \lambda_D\right)^2 - \lambda_K}$$

- lacktriangle Same as eigenvalues of  $\begin{pmatrix} 0 & 1 \\ -\lambda_K & -\lambda_D \end{pmatrix}$
- Can replace K and D (matrices) with corresponding eigenvalues (scalars)
  - Just have to analyze 2x2 system

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# **Split Into More Cases**

- ◆ Still messy! Simplify further
- ◆ If D dominates (e.g. air drag, damping)

$$\alpha \approx \{-\lambda_D, 0\}$$

- Exponential decay and constant
- ◆ If K dominates (e.g. spring force)

$$\alpha \approx \pm \sqrt{-1} \sqrt{\lambda_K}$$

# **Three Test Equations**

- ◆ Constant acceleration (e.g. gravity)
  - a(x,v,t)=g
  - Want exact (2nd order accurate) position
- ◆ Position dependence (e.g. spring force)
  - a(x,v,t)=-Kx
  - Want stability but low or zero damping
  - Look at imaginary axis
- ◆ Velocity dependence (e.g. damping)
  - a(x,v,t)=-Dv
  - Want stability, monotone decay
  - · Look at negative real axis

# **Explicit methods from before**

- ◆ Forward Euler
  - Constant acceleration: bad (1st order)
  - Position dependence: very bad (unstable)
  - Velocity dependence: ok (conditionally monotone/stable)
- ♦ RK3 and RK4
  - Constant acceleration: great (high order)
  - Position dependence: ok (conditionally stable, but damps out oscillation)
  - Velocity dependence: ok (conditionally monotone/stable)

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## Implicit methods from before

- Backward Fuler
  - Constant acceleration: bad (1st order)
  - Position dependence: ok (stable, but damps)
  - Velocity dependence: great (monotone)
- ◆ Trapezoidal Rule
  - Constant acceleration: great (2nd order)
  - Position dependence: great (stable, no damping)
  - Velocity dependence: good (stable but only conditionally monotone)

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# **Setting Up Implicit Solves**

 Let's take a look at actually using Backwards Euler, for example

$$x_{n+1} = x_n + \Delta t \, v_{n+1}$$

$$v_{n+1} = v_n + \Delta t M^{-1} F(x_{n+1}, v_{n+1})$$

◆ Eliminate position, solve for velocity:

$$v_{n+1} = v_n + \Delta t M^{-1} F(x_n + \Delta t v_{n+1}, v_{n+1})$$

♦ Linearize at guess  $v^k$ , solving for  $v_{n+1} \approx v^k + \Delta v$ 

$$v^{k} + \Delta v = v_{n} + \Delta t M^{-1} \left( F\left(x_{n} + \Delta t v^{k}, v^{k}\right) + \Delta t \frac{\partial F}{\partial x} \Delta v + \frac{\partial F}{\partial v} \Delta v \right)$$

◆ Collect terms, multiply by M

$$\left(M - \Delta t \frac{\partial F}{\partial v} - \Delta t^2 \frac{\partial F}{\partial x}\right) \Delta v = M(v_n - v^k) + \Delta t F(x_n + \Delta t v^k, v^k)$$

## **Symmetry**

♦ Why multiply by M?

• Physics often demands that  $\frac{\partial F_{position}}{\partial x}$  and  $\frac{\partial F_{velocity}}{\partial x}$ are symmetric

• And M is symmetric, so this means matrix is symmetric, hence easier to solve

- (physics generally says matrix is SPD even better)
- · If the masses are not equal, the acceleration form of the equations results in an unsymmetric matrix - bad.
- ullet Unfortunately the matrix  $\underline{\partial F_{velocity}}$  is usually unsymmetric
  - Makes solving with it considerably less efficient
  - See Baraff & Witkin, "Large steps in cloth simulation", SIGGRAPH '98 for one solution: throw out bad part

# **Specialized 2nd Order Methods**

- ◆ This is again a big subject
- ◆ Again look at explicit methods, implicit methods
- Also can treat position and velocity dependence differently: mixed implicit-explicit methods

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# **Symplectic Euler**

◆ Like Forward Euler, but updated velocity used for position

$$v_{n+1} = v_n + \Delta t a(x_n, v_n)$$
  
$$x_{n+1} = x_n + \Delta t v_{n+1}$$

- ◆ Some people flip the steps (= relabel v<sub>n</sub>)
- Symplectic means certain qualities of the underlying physics are preserved in discretization - quite desirable visually!
- ◆ [work out test cases]

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# **Symplectic Euler performance**

- ◆ Constant acceleration: bad
  - Velocity right, position off by O(Δt)
- ◆ Position dependence: good
  - Stability limit  $\Delta t < \frac{2}{\sqrt{K}}$
  - No damping! (symplectic)
- ◆ Velocity dependence: ok
  - Monotone limit  $\Delta t < 1/D$
  - Stability limit  $\Delta t < 2/D$

# **Tweaking Symplectic Euler**

- ◆ [sketch algorithms]
- Stagger the velocity to improve x
- ◆ Start off with

$$v_{\frac{1}{2}} = v_0 + \frac{1}{2} \Delta t a(x_0, v_0)$$

◆ Then proceed with

$$v_{n+\frac{1}{2}} = v_{n-\frac{1}{2}} + \frac{1}{2}(t_{n+1} - t_{n-1})a(x_n, v_{n-\frac{1}{2}})$$
  
$$x_{n+1} = x_n + \Delta t v_{n+\frac{1}{2}}$$

◆ Finish off with

$$v_N = v_{N-\frac{1}{2}} + \frac{1}{2} \Delta t a(x_N, v_{N-\frac{1}{2}})$$

# **Staggered Symplectic Euler**

- ◆ Constant acceleration: great!
  - · Position is exact now
- Other cases not effected
  - Was that magic? Main part of algorithm unchanged (apart from relabeling) yet now it's more accurate!
- Only downside to staggering
  - At intermediate times, position and velocity not known together
  - May need to think a bit more about collisions and other interactions with outside algorithms...

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## A common explicit method

◆ May see this one pop up:

$$\begin{aligned} v_{n+1} &= v_n + \Delta t \, a(x_n, v_n) \\ x_{n+1} &= x_n + \Delta t \left(\frac{1}{2} v_n + \frac{1}{2} v_{n+1}\right) = x_n + \Delta t v_n + \frac{1}{2} \Delta t^2 a_n \end{aligned}$$

- ◆ Constant acceleration: great
- ♦ Velocity dependence: ok
  - Conditionally stable/monotone
- ◆ Position dependence: BAD
  - Unconditionally unstable!

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# **An Implicit Compromise**

- Backward Euler is nice due to unconditional monotonicity
  - Although only 1st order accurate, it has the right characteristics for damping
- Trapezoidal Rule is great for everything except damping with large time steps
  - 2nd order accurate, doesn't damp pure oscillation/rotation
- ◆ How can we combine the two?

# **Implicit Compromise**

◆ Use Backward Euler for velocity dependence, Trapezoidal Rule for the rest:

$$x_{n+1} = x_n + \Delta t \left( \frac{1}{2} v_n + \frac{1}{2} v_{n+1} \right)$$

$$v_{n+1} = v_n + \Delta t a \left( \frac{1}{2} x_n + \frac{1}{2} x_{n+1}, v_{n+1}, t_{n+\frac{1}{2}} \right)$$

- ◆ Constant acceleration: great (2nd order)
- Position dependence: great (2nd order, no damping)
- Velocity dependence: great (unconditionally monotone)

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