

## Notes

- ◆ More reading on web site
  - Baraff & Witkin's classic cloth paper
  - Grinspun et al. on bending
  - Optional: Teran et al. on FVM in graphics

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## Hyper-elasticity

- ◆ Another common way to look at elasticity
  - Useful for handling weird nonlinear compressibility laws, for reduced dimension models, and more
- ◆ Instead of defining stress, define an elastic potential energy
  - Strain energy density  $W=W(A)$
  - $W=0$  for no deformation,  $W>0$  for deformation
  - Total potential energy is integral of  $W$  over object
- ◆ This is called hyper-elasticity or Green elasticity
- ◆ For most (the ones that make sense) stress-strain relationships can define  $W$ 
  - E.g. linear relationship:  $W=\sigma:\epsilon=\text{trace}(\sigma^T\epsilon)$

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## Variational Derivatives

- ◆ Force is the negative gradient of potential
  - Just like gravity
- ◆ What does this mean for a continuum?
  - $W=W(\partial X/\partial p)$ , how do you do  $-d/dX$ ?
- ◆ Variational derivative:  $W_{total}[X + \epsilon Y] = \int W \left( \frac{\partial X}{\partial p} + \epsilon \frac{\partial Y}{\partial p} \right)$ 
  - So variational derivative is  $-\nabla \cdot \partial W / \partial A$
  - And  $f = \nabla \cdot \partial W / \partial A$
  - Then stress is  $\partial W / \partial A$

$$\begin{aligned} &\approx \int W \left( \frac{\partial X}{\partial p} \right) + \epsilon \frac{\partial W}{\partial A} \frac{\partial Y}{\partial p} \\ &= W_{total} + \epsilon \int \frac{\partial W}{\partial A} \frac{\partial Y}{\partial p} \\ &= W_{total} - \epsilon \int Y \nabla \cdot \frac{\partial W}{\partial A} \end{aligned}$$

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## Numerics

- ◆ Simpler approach: find discrete  $W_{total}$  as a sum of  $W$ 's for each element
  - Evaluate just like FEM, or any way you want
- ◆ Take gradient w.r.t. positions  $\{x_i\}$ 
  - Ends up being a Galerkin method
- ◆ Also note that an implicit method might need Jacobian = negative Hessian of energy
  - Must be symmetric, and at least near stable configurations must be negative definite

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## Curve / Springs

- ◆ Take  $W(A) = 1/2 E(|A|-1)^2 L$  for each segment
  - Note factor of  $L$ : this is approximation to an integral over segment in object space of length  $L$
- ◆  $A = (x_{i+1} - x_i) / L$  is the deformation gradient for piecewise linear elements
- ◆ Then take derivative w.r.t.  $x_i$  to get this element's contribution to force on  $i$
- ◆ Lo and behold [exercise] get exactly the original spring force
- ◆ Note: defining stress and strain would be more complicated, because of the dimension differences
  - $A$  is  $3 \times 1$ , not square

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## Surface elasticity

- ◆ For linear stress-strain, can use  $W(A) = \sigma : G = \sigma_{ij} G_{ij}$
- ∪ The simplest model from before gives  $W = \lambda G_{kk}^2 + \mu G_{ij} G_{ij}$
- ∪ Remember  $G = 1/2(A^T A - I)$
- ∪ Tedious to differentiate, but doable
  - Tensors and chain rule over and over
- ∪ Let's leave it that
  - In practice, springs with speed-of-sound heuristic are good enough most of the time

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## Bending energy

- ◆ Bending is very difficult to get a handle on without variational approach
- ◆ Bending strain energy density:  
 $W = 1/2 B \kappa^2$
- ∪ Here  $\kappa$  is mean curvature
  - Look at circles that fit surface
  - Maximum radius R and minimum radius r
  - $\kappa = (1/R + 1/r)/2$
  - Can define directly from second derivatives of  $X(p)$
  - Uh-oh - second derivatives? [FEM nastier]
  - W is 2nd order, stress is 3rd order, force is 4th order derivatives!

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## Fourth order problems

- ◆ Linearize and simplify drastically, look for steady-state solution ( $F=0$ ): spline equations
  - Essentially 4th derivatives are zero
  - Solutions are (bi-)cubics
- ◆ Model (nonsteady) problem:  $X_{tt} = -X_{pppp}$ 
  - Assume solution  $x(p, t) = e^{\sqrt{-ik}(p-ct)}$
  - Wave of spatial frequency k, moving at speed c
  - [solve for wave parameters]
  - Dispersion relation: small waves move really fast
  - CFL limit (and stability): for fine grids, BAD
  - Thankfully, we rarely get that fine

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## Implicit/Explicit Methods

- ◆ Implicit bending is painful
- ◆ In graphics, usually unnecessary
  - Dominant forces on the grid resolution we use tend to be the 2nd order terms: stretching etc.
- ◆ But nice to go implicit to avoid time step restriction for stretching terms
- ◆ No problem: treat some terms (bending) explicitly, others (stretching) implicitly
  - $v_{n+1} = v_n + \Delta t/m(F_1(x_n, v_n) + F_2(x_{n+1}, v_{n+1}))$
  - All bending is in  $F_1$ , half the elastic stretch in  $F_1$ , half the elastic stretch in  $F_2$ , all the damping in  $F_2$

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## Discrete Mean Curvature

- ◆ [draw triangle pair]  $\kappa \sim \frac{\theta}{h_1 + h_2}$
- ∪  $\kappa$  for that chunk varies as
- ◆ So integral of  $\kappa^2$  varies as  $W = \sum_e \frac{\theta^2}{(h_1 + h_2)^2} (|\Delta_1| + |\Delta_2|)$
- $\sim \sum_e \frac{\theta^2 |e|^2}{|\Delta_1| + |\Delta_2|}$
- ◆ Edge length, triangle areas, normals are all easy to calculate
- ∪  $\theta$  needs inverse trig functions
- ∪ But  $\theta^2$  behaves a lot like  $1 - \cos(\theta/2)$  over interval  $[-\pi, \pi]$  [draw picture]

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